

On the field recursion method for two-component composites

Aaron Welters

Department of Mathematical Sciences
Florida Institute of Technology
Melbourne, FL USA

Joint work with:

Maxence Cassier (Institut Fresnel) & Graeme Milton (Univ. of Utah)

2018

1. G. W. Milton (editor), *Extending the Theory of Composites to Other Areas of Science*, Milton-Patton Publishing, 2016. See book review in *SIAM Review*, Vol. 60, No. 2, June 2018.
2. M. Cassier, A. Welters, and G. W. Milton, *A rigorous approach to the field recursion method for two-component composites with isotropic phases*, Chap. 10.
3. G. W. Milton, *The Theory of Composites*, Cambridge University Press, 2002. See, in particular, Chap. 29.
4. G. W. Milton, *Multicomponent composites, electrical networks, and new types of continued fractions. I, II*, *Comm. Math. Phys.* 111 (2)&(3): 281–327 & 329–372, 1987.
5. G. W. Milton, *The field equation recursion method*, in G. Dal Maso and G. F. Dell'Antonio (eds.), *Composite Media and Homogenization Theory*, Vol. 5 of *Progress in Nonlinear Diff. Eqs. & Their Apps.*, pp. 224–245, Birkhauser, 1991.
6. G. W. Milton and K. Golden, *Thermal conduction in composites*, in T. Ashworth and D. R. Smith (eds.), *Thermal Conductivity* 18, pp. 571–582, Plenum, 1985.

Motivation

- For two-component composites, the bounds calculated from the field equation recursion method, variational principles, or the analytic method.
- **The real power of the recursion method:** multiphase composites – Obtain bounds on complex effective moduli where other methods fail & many of the bounds are sharp.
- Advantage: applicable to any problem formulizable in the abstract theory of composites and **it produces bounds in a systematic way.**
- Two-component composite with isotropic phase: (i) wanted a **rigorous derivation** of the recursion method **using the abstract theory of composites.**
- (ii) Wanted to **significantly reduce**, at each level of induction, the **necessary assumptions** required to proceed to the next level.
- (iii) Wanted a **simplification** of many of the **mathematical aspects** of the method.

Hilbert space orthogonal-triple, Z-problem, & effective operator

- **Hilbert space:** \mathcal{H} ; inner product: (\cdot, \cdot) .
- **Orthogonal-triple decomposition:** $\mathcal{H} = \mathcal{U} \oplus \mathcal{E} \oplus \mathcal{J}$.
- **Constitutive relations:** For a given bounded linear operator $L \in \mathcal{L}(\mathcal{H})$,

$$J = LE, \text{ where } E \in \mathcal{U} \oplus \mathcal{E}, J \in \mathcal{U} \oplus \mathcal{J}.$$

- **Z-problem:** Given $E_0 \in \mathcal{U}$, find $(J_0, E, J) \in \mathcal{U} \times \mathcal{E} \times \mathcal{J}$ such that
- $$J_0 + J = L(E_0 + E).$$

- **Effective operator:** If there exists an $Z_* \in \mathcal{L}(\mathcal{U})$ such that

$$J_0 = L_* E_0$$

for every $E_0 \in \mathcal{U}$, where (J_0, E, J) is a solution of the Z-problem associated with E_0 , then we call L_* an effective operator.

Z-problem solution & effective operator as a Schur complement

- Express L as a 3×3 block operator matrix with respect to

$$\mathcal{H} = \mathcal{U} \oplus \mathcal{E} \oplus \mathcal{J}:$$

$$L = \begin{bmatrix} L_{00} & L_{01} & L_{02} \\ L_{10} & L_{11} & L_{12} \\ L_{20} & L_{21} & L_{22} \end{bmatrix}, \quad L_{\mathcal{U} \oplus \mathcal{E}} := \begin{bmatrix} L_{00} & L_{01} \\ L_{10} & L_{11} \end{bmatrix}, \quad L_{\mathcal{E}} := L_{11}.$$

- The Z-problem for a given $E_0 \in \mathcal{U}$, equivalent to solving:

$$\begin{cases} J_0 = L_{00}E_0 + L_{01}E, \\ 0 = L_{10}E_0 + L_{11}E, \\ J = L_{20}E_0 + L_{21}E. \end{cases}$$

for $(J_0, E, J) \in \mathcal{U} \times \mathcal{E} \times \mathcal{J}$.

- If L_{11}^{-1} exists then a **unique solution exists to the Z-problem**, the **effective operator L_* exists and is unique**, and

$$J = L_{20}E_0 + L_{21}E, \quad E = -L_{11}^{-1}L_{10}E_0, \quad J_0 = L_*E_0, \quad (1)$$

$$L_* = L_{\mathcal{U} \oplus \mathcal{E}} / L_{\mathcal{E}} = L_{00} - L_{01}L_{11}^{-1}L_{10} \quad (\text{Schur complement}) \quad (2)$$

Orthogonal $Z(2)$ subspace collection

- Orthogonal $Z(2)$ subspace collection: $\mathcal{H}, \mathcal{U}, \mathcal{E}, \mathcal{J}, \mathcal{P}_j, j = 1, 2$, where

$$\mathcal{H} = \mathcal{U} \overset{\perp}{\oplus} \mathcal{E} \overset{\perp}{\oplus} \mathcal{J} = \mathcal{P}_1 \overset{\perp}{\oplus} \mathcal{P}_2.$$

- Orthogonal projections: $\Gamma_0, \Gamma_1, \Gamma_2, \Lambda_j, j = 1, 2$, where

$$\Gamma_0 \mathcal{H} = \mathcal{U}, \Gamma_1 \mathcal{H} = \mathcal{E}, \Gamma_2 \mathcal{H} = \mathcal{J}, \Lambda_j \mathcal{H} = \mathcal{P}_j.$$

- Two-component composite:

$$L = L(L_1, L_2) = L_1 \Lambda_1 + L_2 \Lambda_2,$$

where $L_1, L_2 \in \mathcal{L}(\mathcal{H})$.

- Isotropic, two-phase:

$$L = L(L_1, L_2) = L_1 \Lambda_1 + L_2 \Lambda_2, \quad L_1, L_2 \in \mathbb{C}.$$

Properties: Homogeneity, Normalization, & Herglotz

- **Schur complement formula (2) implies** homogeneity, normalization, & Herglotz properties.
- **Homogeneity** property:

$$L_*(cL_1, cL_2) = cL_*(L_1, L_2), \quad \forall c \in \mathbb{C} \setminus \{0\}.$$

- **Normalization** property:

$$L_*(1, 1) = I_{\mathcal{U}}.$$

- **Herglotz** property: The function $L_* : (\mathbb{C}^+)^2 \rightarrow \mathcal{L}(\mathcal{U})$ is analytic and

$$\operatorname{Im} L_*(L_1, L_2) > 0, \quad \forall (L_1, L_2) \in (\mathbb{C}^+)^2,$$

where

$$\mathbb{C}^+ = \{z \in \mathbb{C} : \operatorname{Im} z > 0\}, \quad \operatorname{Im} M = \frac{M - M^*}{2i}, \quad M \in \mathcal{L}(\mathcal{H}).$$

Quintessential example: conductivity equation

Framework for **conductivity equation** ($\nabla \cdot \sigma \nabla u = 0$) of a **two-component periodic composite** with unit cell $\mathcal{D} = [0, 2\pi]^d$ ($d = 2, 3$).

- **Hilbert space:** $\mathcal{H} = [L_{\text{per}}^2(\mathcal{D})]^d$; inner product and average
$$(E, F) = \frac{1}{|\mathcal{D}|} \int_{\mathcal{D}} E(x)^T \overline{F(x)} dx, \quad \langle F \rangle = \frac{1}{|\mathcal{D}|} \int_{\mathcal{D}} F(x) dx, \quad \forall E, F \in \mathcal{H}.$$
- **Orthogonal-triple** (Hodge) decomposition: $\mathcal{H} = \mathcal{U} \oplus \mathcal{E} \oplus \mathcal{J}$, where
$$\begin{aligned} \mathcal{U} &= \left\{ U \in \mathcal{H} \mid U \equiv C, \text{ for some } C \in \mathbb{C}^d \right\}; \\ \mathcal{E} &= \left\{ E \in \mathcal{H} \mid \nabla \times E = 0 \text{ and } \langle E \rangle = 0 \right\}; \\ \mathcal{J} &= \left\{ J \in \mathcal{H} \mid \nabla \cdot J = 0 \text{ and } \langle J \rangle = 0 \right\}. \end{aligned}$$
- **Constitutive relation:** Multiplication operator
$$\sigma = \sigma(x) \in M_d(L_{\text{per}}^\infty(\mathcal{D})),$$
$$J = \sigma E, \text{ where } E \in \mathcal{U} \oplus \mathcal{E}, J \in \mathcal{U} \oplus \mathcal{J}.$$
- **Z-problem:** Given $E_0 \in \mathcal{U}$, find $(J_0, E, J) \in \mathcal{U} \times \mathcal{E} \times \mathcal{J}$ such that
$$J_0 + J = \sigma(E_0 + E).$$
- **Effective operator** $\sigma_* \in \mathcal{L}(\mathcal{U})$:
$$J_0 = \sigma_* E_0 \text{ or, equivalently, } \langle J_0 + J \rangle = \sigma_* \langle E_0 + E \rangle.$$

Quintessential example: conductivity equation

- Two-component composite, isotropic, two-phase:

$$\mathcal{D} = \mathcal{D}_1 \cup \mathcal{D}_2, \quad \mathcal{D}_1 \cap \mathcal{D}_2 = \emptyset, \quad \chi_j \in L^2_{\text{per}}(\mathcal{D})$$

$$\text{for } x \in \mathcal{D}_j, \quad \chi_j(x) = \begin{cases} 1 & \text{if } x \in \mathcal{D}_j, \\ 0 & \text{if } x \notin \mathcal{D}_j. \end{cases} \quad (\text{characteristic function of } \mathcal{D}_j)$$

$$\sigma = \sigma(\sigma_1, \sigma_2) = \sigma_1 \chi_1 I + \sigma_2 \chi_2 I, \quad \sigma_1, \sigma_2 \in \mathbb{C}.$$

- Orthogonal $Z(2)$ subspace collection:

$$\mathcal{H} = \mathcal{U} \oplus \mathcal{E} \oplus \mathcal{J} = \mathcal{P}_1 \oplus \mathcal{P}_2.$$

- Orthogonal projections: $\Gamma_0, \Gamma_1, \Gamma_2, \Lambda_j, j = 1, 2$, where

$$\Gamma_0 \mathcal{H} = \mathcal{U}, \quad \Gamma_1 \mathcal{H} = \mathcal{E}, \quad \Gamma_2 \mathcal{H} = \mathcal{J}, \quad \Lambda_j \mathcal{H} = \mathcal{P}_j.$$

- For any $F \in \mathcal{H} = [L^2_{\text{per}}(\mathcal{D})]^d$,

$$\Gamma_0 F = \langle F \rangle, \quad \Gamma_1 F = [\nabla(\nabla \cdot \nabla)^{-1} \nabla \cdot] F, \quad \Gamma_2 F = [\nabla \times (-\nabla \cdot \nabla)^{-1} \nabla \times] F, \\ \Lambda_j F = \chi_j F.$$

Quintessential example: conductivity equation

- Explicitly, for any $F \in \mathcal{H} = [L^2_{\text{per}}(\mathcal{D})]^d$,

$$F(x) = \widehat{F}(0) + \sum_{0 \neq k \in \mathbb{Z}^d} \widehat{F}(k) e^{ik \cdot x}, \quad I = \Gamma_0 + \Gamma_1 + \Gamma_2,$$

$$\Gamma_0 F = \langle F \rangle = \widehat{F}(0),$$

$$\Gamma_1 F = [\nabla(\nabla \cdot \nabla)^{-1} \nabla \cdot] F = \sum_{0 \neq k \in \mathbb{Z}^d} \frac{k}{|k|^2} k \cdot \widehat{F}(k) e^{ik \cdot x},$$

$$\begin{aligned} \Gamma_2 F &= [\nabla \times (-\nabla \cdot \nabla)^{-1} \nabla \times] F = \sum_{0 \neq k \in \mathbb{Z}^d} \frac{k \times (k \times \widehat{F}(k))}{|k|^2} e^{ik \cdot x} \\ &= \sum_{0 \neq k \in \mathbb{Z}^d} \left(1 - \frac{k}{|k|^2} k \cdot \right) \widehat{F}(k) e^{ik \cdot x}. \end{aligned}$$

- There are **other examples**, e.g., **Dirichlet-to-Neumann map**, which also fits into the abstract theory of composites **as an effective operator**. See [1]!

The Y -problem and Y -operator

- **Hilbert space:** \mathcal{K} ; inner product: (\cdot, \cdot) .
- **2-orthogonal decompositions:** $\mathcal{K} = \mathcal{E} \oplus \mathcal{J} = \mathcal{V} \oplus \mathcal{H}$.
- **Y -problem:** Given $L \in \mathcal{L}(\mathcal{H})$ and $E_1 \in \mathcal{V}$, find $(E, J) \in \mathcal{E} \times \mathcal{J}$ such that there exists $J_1 \in \mathcal{V}$, $E_2, J_2 \in \mathcal{H}$ satisfying

$$E = E_1 + E_2, \quad J = J_1 + J_2, \quad \text{and} \quad J_2 = LE_2.$$

- **Y -operator:** If there exists an $Y_* \in \mathcal{L}(\mathcal{V})$ such that

$$J_1 = -Y_* E_1$$

for every $E_1 \in \mathcal{V}$, whenever (E, J) is a solution of the Y -problem associated with E_1 , then we call Y_* a Y -operator.

- Orthogonal projections: $\Gamma_j, \Pi_j, j = 1, 2$, where

$$\Gamma_1 \mathcal{K} = \mathcal{E}, \Gamma_2 \mathcal{K} = \mathcal{J}, \Pi_1 \mathcal{K} = \mathcal{V}, \Pi_2 \mathcal{K} = \mathcal{H}.$$

- i.e., **Y-problem**: Given $L \in \mathcal{L}(\mathcal{H})$ and $E_1 \in \mathcal{V}$, find $(E, J) \in \mathcal{E} \times \mathcal{J}$ such that $\Pi_2 J = L \Pi_2 E$ and $\Pi_1 E = E_1$.

Proposition (Y-Problem: Existence & Uniqueness)

If $L^{-1} : \mathcal{H} \rightarrow \mathcal{H}$ and $(\Gamma_2 \Pi_2 L^{-1} \Pi_2 \Gamma_2)^{-1} : \mathcal{J} \rightarrow \mathcal{J}$ exist then (i) solution is **unique**; (ii) **exists** with hypothesis $\Pi_2 L^{-1} \Pi_2 J = L^{-1} \Pi_2 J$:

$$J = -(\Gamma_2 \Pi_2 L^{-1} \Pi_2 \Gamma_2)^{-1} \Gamma_2 E_1, \quad E = E_1 + E_2, \quad E_2 = L^{-1} \Pi_2 J, \quad (3)$$

$$Y_* = \Pi_1 \Gamma_2 (\Gamma_2 \Pi_2 L^{-1} \Pi_2 \Gamma_2)^{-1} \Gamma_2 \Pi_1. \quad (4)$$

Proof.

$$\begin{aligned} \text{(i)} \quad \Pi_2 J = L \Pi_2 E, \Pi_1 E = E_1 &\Rightarrow 0 = \Gamma_2 E = \Gamma_2 \Pi_1 E + \Gamma_2 \Pi_2 E \\ &= \Gamma_2 E_1 + \Gamma_2 L^{-1} \Pi_2 J = \Gamma_2 E_1 + \Gamma_2 \Pi_2 L^{-1} \Pi_2 \Gamma_2 J; \end{aligned}$$

$$\begin{aligned} \text{(ii)} \quad (3) &\Rightarrow \Pi_j E = E_j, \Pi_2 J = L \Pi_2 E, J = \Gamma_2 J, \Gamma_2 E = \Gamma_2 E_1 + \Gamma_2 E_2 \\ &= \Gamma_2 E_1 + \Gamma_2 L^{-1} \Pi_2 J = \Gamma_2 E_1 + (\Gamma_2 \Pi_2 L^{-1} \Pi_2 \Gamma_2) J = 0 \Rightarrow (E, J) \in \mathcal{E} \times \mathcal{J} \end{aligned}$$

Orthogonal $Y(2)$ subspace collection

- Orthogonal $Y(2)$ subspace collection: $\mathcal{K}, \mathcal{E}, \mathcal{J}, \mathcal{V}, \mathcal{H}, \mathcal{P}_j, j = 1, 2$, where

$$\mathcal{K} = \mathcal{E} \overset{\perp}{\oplus} \mathcal{J} = \mathcal{V} \overset{\perp}{\oplus} \mathcal{H}, \quad \mathcal{H} = \mathcal{P}_1 \overset{\perp}{\oplus} \mathcal{P}_2.$$

- Orthogonal projections: $\Lambda_j : \mathcal{H} \rightarrow \mathcal{H}, j = 1, 2$, where

$$\Lambda_j \mathcal{H} = \mathcal{P}_j.$$

- Two-component composite:

$$L = L(L_1, L_2) = L_1 \Lambda_1 + L_2 \Lambda_2,$$

where $L_1, L_2 \in \mathcal{L}(\mathcal{H})$.

- Isotropic, two-phase:

$$L = L(L_1, L_2) = L_1 \Lambda_1 + L_2 \Lambda_2, \quad L_1, L_2 \in \mathbb{C}.$$

Properties: Homogeneity, Positivity, & Herglotz

- If $\mathcal{V} \cap \mathcal{J} = \emptyset$ then **formula (4) implies** homogeneity, positivity, & Herglotz properties.
- **Homogeneity** property:

$$Y_*(cL_1, cL_2) = cY_*(L_1, L_2), \quad \forall c \in \mathbb{C} \setminus \{0\}.$$

- **Positivity** property:

$$Y_*(1, 1) = \Pi_1 \Gamma_2 (\Gamma_2 \Pi_2 \Gamma_2)^{-1} \Gamma_2 \Pi_1 |_{\mathcal{V}} > 0.$$

- **Herglotz** property: The function $Y_* : (\mathbb{C}^+)^2 \rightarrow \mathcal{L}(\mathcal{V})$ is analytic and

$$\operatorname{Im} Y_*(L_1, L_2) > 0, \quad \forall (L_1, L_2) \in (\mathbb{C}^+)^2.$$

Hint: Use identity $\operatorname{Im} (M^{-1}) = -M^{-1} (\operatorname{Im} M) (M^{-1})^*$.

The field recursion method: initial conditions & base case

Field equation recursion method: isotropic, two-phase

(!) Indicates continued fraction formulas

Base case $i = 0$:

- **Input** – orthogonal $Z(2)$ subspace collection:

$$\mathcal{H}^{(0)}, \mathcal{U}^{(0)}, \mathcal{E}^{(0)}, \mathcal{J}^{(0)}, \mathcal{P}_j^{(0)}, j = 1, 2.$$

(1) Generate – the effective operator $L_*^{(0)}$;

(2) Generate – orthogonal $Y(2)$ subspace collection:

$$\mathcal{K}^{(0)}, \mathcal{E}^{(0)}, \mathcal{J}^{(0)}, \mathcal{V}^{(0)}, \mathcal{H}^{(1)}, \mathcal{P}_j^{(1)}, j = 1, 2; \text{ the } Y\text{-operator } Y_*^{(0)};$$

(3) Generate – orthogonal $Z(2)$ subspace collection:

$$\mathcal{H}^{(1)}, \mathcal{U}^{(1)}, \mathcal{E}^{(1)}, \mathcal{J}^{(1)}, \mathcal{P}_j^{(1)}, j = 1, 2; \text{ the } Z\text{-operator } L_*^{(1)}.$$

(4) Formulas, properties, & relations:

$$L^{(0)} = L, (!) (L_*^{(0)}, Y_*^{(0)}, L_*^{(1)}).$$

- **Output** – the new $Z(2)$ subspace collection:

$$\mathcal{H}^{(1)}, \mathcal{U}^{(1)}, \mathcal{E}^{(1)}, \mathcal{J}^{(1)}, \mathcal{P}_j^{(1)}, j = 1, 2.$$

The field recursion method: Induction step

Field equation recursion method: isotropic, two-phase

(!) Indicates continued fraction formulas

Induction step $i, i \geq 1$:

- **Input** – orthogonal $Z(2)$ subspace collections:

$$\mathcal{H}^{(i)}, \mathcal{U}^{(i)}, \mathcal{E}^{(i)}, \mathcal{J}^{(i)}, \mathcal{P}_j^{(i)}, j = 1, 2.$$

(1) Generate – the effective operator $L_*^{(i)}$;

(2) Generate – orthogonal $Y(2)$ subspace collection:

$$\mathcal{K}^{(i)}, \mathcal{E}^{(i)}, \mathcal{J}^{(i)}, \mathcal{V}^{(i)}, \mathcal{H}^{(i+1)}, \mathcal{P}_j^{(i+1)}, j = 1, 2; \text{ the } Y\text{-operator } Y_*^{(i)};$$

(3) Generate – orthogonal $Z(2)$ subspace collection:

$$\mathcal{H}^{(i+1)}, \mathcal{U}^{(i+1)}, \mathcal{E}^{(i+1)}, \mathcal{J}^{(i+1)}, \mathcal{P}_j^{(i+1)}, j = 1, 2; \text{ the } Z\text{-operator } L_*^{(i+1)}.$$

(4) Formulas, properties, & relations:

$$L^{(0)} = L, (!) \left(L_*^{(i)}, Y_*^{(i)}, L_*^{(i+1)} \right).$$

- **Output** – the new $Z(2)$ subspace collection:

$$\mathcal{H}^{(i+1)}, \mathcal{U}^{(i+1)}, \mathcal{E}^{(i+1)}, \mathcal{J}^{(i+1)}, \mathcal{P}_j^{(i+1)}, j = 1, 2.$$

- Input – orthogonal $Z(2)$ subspace collection:

$$\mathcal{H}^{(0)}, \mathcal{U}^{(0)}, \mathcal{E}^{(0)}, \mathcal{J}^{(0)}, \mathcal{P}_j^{(0)}, j = 1, 2.$$

- (2) Generate – orthogonal $Y(2)$ subspace collection

$$\mathcal{K}^{(0)}, \mathcal{E}^{(0)}, \mathcal{J}^{(0)}, \mathcal{V}^{(0)}, \mathcal{H}^{(1)}, \mathcal{P}_j^{(1)}, j = 1, 2:$$

$$\mathcal{K}^{(0)} = \mathcal{H}^{(0)} \ominus \mathcal{U}^{(0)}, \mathcal{P}_j^{(1)} = \mathcal{P}_j^{(0)} \cap \mathcal{K}^{(0)}, j = 1, 2,$$

$$\mathcal{H}^{(1)} = \mathcal{P}_1^{(1)} \oplus \mathcal{P}_2^{(1)}, \mathcal{V}^{(0)} = \mathcal{K}^{(0)} \ominus \mathcal{H}^{(1)}.$$

- (3) Generate – orthogonal $Z(2)$ subspace collection

$$\mathcal{H}^{(1)}, \mathcal{U}^{(1)}, \mathcal{E}^{(1)}, \mathcal{J}^{(1)}, \mathcal{P}_j^{(1)}, j = 1, 2:$$

$$\mathcal{E}^{(1)} = \mathcal{E}^{(0)} \cap \mathcal{H}^{(1)}, \mathcal{J}^{(1)} = \mathcal{J}^{(0)} \cap \mathcal{H}^{(1)}, \mathcal{U}^{(1)} = \mathcal{H}^{(1)} \ominus \left[\mathcal{E}^{(1)} \oplus \mathcal{J}^{(1)} \right]$$

- (4) Formulas, properties, & relations:

$$\begin{aligned} L_*^{(l)} &= \Gamma_0^{(l)} L \Gamma_0^{(l)} - \Gamma_0^{(l)} L \Gamma_1^{(l)} \left(\Gamma_1^{(l)} L \Gamma_1^{(l)} \right)^{-1} \Gamma_1^{(l)} L \Gamma_0^{(l)} \\ &= \Gamma_0^{(l)} \left[\left(\Gamma_0^{(l)} + \Gamma_2^{(l)} \right) L^{-1} \left(\Gamma_0^{(l)} + \Gamma_2^{(l)} \right) \right]^{-1} \Gamma_0^{(l)}, l = 0, 1 \end{aligned}$$

$$(!) Y_*^{(0)} = \Pi_1^{(0)} \Gamma_2^{(0)} \left(\Gamma_2^{(0)} \Pi_2^{(0)} L^{-1} \Pi_2^{(0)} \Gamma_2^{(0)} \right)^{-1} \Gamma_2^{(0)} \Pi_1^{(0)} = K^{(0)} L_*^{(1)} \left(K^{(0)} \right)^*$$

$$K^{(0)} \in \mathcal{L} \left(\mathcal{U}^{(1)}, \mathcal{V}^{(0)} \right) \text{ invertible, } K^{(0)} = - \left(\Pi_1^{(0)} \Gamma_1^{(0)} \Pi_1^{(0)} \right)^{-1} \Pi_1^{(0)} \Gamma_1^{(0)} \Pi_2^{(0)},$$

$$(!) L_*^{(0)} = \Gamma_0^{(0)} L \Gamma_0^{(0)} - \Gamma_0^{(0)} L \Pi_1^{(0)} \left(\Pi_1^{(0)} L \Pi_1^{(0)} + Y_*^{(0)} \right)^{-1} \Pi_1^{(0)} L \Gamma_0^{(0)}.$$

Quintessential example revisited: conductivity equation

$$\nabla \cdot \sigma \nabla u = 0, \text{ unit cell } \mathcal{D} = [0, 2\pi]^d = \mathcal{D}_1 \cup \mathcal{D}_2, \mathcal{H}^{(0)} = [L_{\text{per}}^2(\mathcal{D})]^d,$$

$$\mathcal{U}^{(0)} = \left\{ U \in \mathcal{H}^{(0)} \mid U \equiv C, \text{ for some } C \in \mathbb{C}^d \right\};$$

$$\mathcal{E}^{(0)} = \left\{ E \in \mathcal{H}^{(0)} \mid \nabla \times E = 0 \text{ and } \langle E \rangle = 0 \right\};$$

$$\mathcal{J}^{(0)} = \left\{ J \in \mathcal{H}^{(0)} \mid \nabla \cdot J = 0 \text{ and } \langle J \rangle = 0 \right\}.$$

$$\mathcal{H}^{(1)} = \mathcal{U}^{(1)} \dot{\oplus} \mathcal{E}^{(1)} \dot{\oplus} \mathcal{J}^{(1)} = \mathcal{P}_1^{(1)} \dot{\oplus} \mathcal{P}_2^{(1)}, \mathcal{K}^{(0)} = \mathcal{E}^{(0)} \dot{\oplus} \mathcal{J}^{(0)} = \mathcal{V}^{(0)} \dot{\oplus} \mathcal{H}^{(1)}.$$

$$\mathcal{P}_j^{(1)} = \mathcal{P}_j^{(0)} \cap \mathcal{K}^{(0)} = \left\{ F \in \mathcal{H}^{(0)} : \chi_j F = F \text{ and } \langle F \rangle = 0 \right\}, j = 1, 2.$$

$$\mathcal{H}^{(1)} = \left\{ F \in \mathcal{H}^{(0)} : \langle \chi_j F \rangle = 0, j = 1, 2 \right\}.$$

$$\mathcal{V}^{(0)} = \mathcal{K}^{(0)} \dot{\ominus} \mathcal{H}^{(1)} = \left\{ F \in \mathcal{H}^{(0)} : \chi_j F|_{\mathcal{D}_j} \equiv C_j \in \mathbb{C}^d, j = 1, 2 \text{ and } \langle F \rangle = 0 \right\}$$

$$\mathcal{E}^{(1)} = \left\{ E \in \mathcal{H}^{(1)} : \nabla \times E = 0 \right\}, \mathcal{J}^{(1)} = \left\{ J \in \mathcal{H}^{(1)} : \nabla \cdot J = 0 \right\}.$$

$$\mathcal{U}^{(1)} = \mathcal{H}^{(1)} \dot{\ominus} \left[\mathcal{E}^{(1)} \dot{\oplus} \mathcal{J}^{(1)} \right] \text{ (more complicated).}$$

- Input – for $i \geq 1$, orthogonal $Z(2)$ subspace collection:

$$\mathcal{H}^{(i)}, \mathcal{U}^{(i)}, \mathcal{E}^{(i)}, \mathcal{J}^{(i)}, \mathcal{P}_j^{(i)}, j = 1, 2.$$

- (2) Generate – orthogonal $Y(2)$ subspace collection

$$\mathcal{K}^{(i)}, \mathcal{E}^{(i)}, \mathcal{J}^{(i)}, \mathcal{V}^{(i)}, \mathcal{H}^{(i+1)}, \mathcal{P}_j^{(i+1)}, j = 1, 2:$$

$$\mathcal{K}^{(i)} = \mathcal{H}^{(i)} \ominus \mathcal{U}^{(i)}, \mathcal{P}_j^{(i+1)} = \mathcal{P}_j^{(i)} \cap \mathcal{K}^{(i)}, j = 1, 2,$$

$$\mathcal{H}^{(1)} = \mathcal{P}_1^{(i+1)} \oplus \mathcal{P}_2^{(i+1)}, \mathcal{V}^{(i)} = \mathcal{K}^{(i)} \ominus \mathcal{H}^{(i+1)}.$$

- (3) Generate – orthogonal $Z(2)$ subspace collection

$$\mathcal{H}^{(i+1)}, \mathcal{U}^{(i+1)}, \mathcal{E}^{(i+1)}, \mathcal{J}^{(i+1)}, \mathcal{P}_j^{(i+1)}, j = 1, 2:$$

$$\mathcal{E}^{(i+1)} = \mathcal{E}^{(i)} \cap \mathcal{H}^{(i+1)}, \mathcal{J}^{(i+1)} = \mathcal{J}^{(i)} \cap \mathcal{H}^{(i+1)}, \mathcal{U}^{(i+1)} = \mathcal{H}^{(i+1)} \ominus [\mathcal{E}^{(i+1)} \oplus \mathcal{J}^{(i+1)}]$$

- (4) Formulas, properties, & relations:

$$L_*^{(i+l)} = \Gamma_0^{(i+l)} L \Gamma_0^{(i+l)} - \Gamma_0^{(i+l)} L \Gamma_1^{(i+l)} \left(\Gamma_1^{(i+l)} L \Gamma_1^{(i+l)} \right)^{-1} \Gamma_1^{(i+l)} L \Gamma_0^{(i+l)}$$

$$= \Gamma_0^{(i+l)} \left[\left(\Gamma_0^{(i+l)} + \Gamma_2^{(i+l)} \right) L^{-1} \left(\Gamma_0^{(i+l)} + \Gamma_2^{(i+l)} \right) \right]^{-1} \Gamma_0^{(i+l)}, l = 0, 1$$

$$(!) Y_*^{(i)} = \Pi_1^{(i)} \Gamma_2^{(i)} \left(\Gamma_2^{(i)} \Pi_2^{(i)} L^{-1} \Pi_2^{(i)} \Gamma_2^{(i)} \right)^{-1} \Gamma_2^{(i)} \Pi_1^{(i)} = K^{(i)} L_*^{(i+1)} \left(K^{(i)} \right)^*,$$

$$K^{(i)} \in \mathcal{L} \left(\mathcal{U}^{(i+1)}, \mathcal{V}^{(i)} \right) \text{ invertible, } K^{(i)} = - \left(\Pi_1^{(i)} \Gamma_1^{(i)} \Pi_1^{(i)} \right)^{-1} \Pi_1^{(i)} \Gamma_1^{(i)} \Pi_2^{(i)},$$

$$(!) L_*^{(i)} = \Gamma_0^{(i)} L \Gamma_0^{(i)} - \Gamma_0^{(i)} L \Pi_1^{(i)} \left(\Pi_1^{(i)} L \Pi_1^{(i)} + Y_*^{(i)} \right)^{-1} \Pi_1^{(i)} L \Gamma_0^{(i)}.$$

Theorem ([2], Cassier-Milton-Welters)

Assume $\mathcal{H}^{(0)}, \mathcal{U}^{(0)}, \mathcal{E}^{(0)}, \mathcal{J}^{(0)}, \mathcal{P}_j^{(0)}$, $j = 1, 2$ is an orthogonal $Z(2)$ subspace collection, $\mathcal{U}^{(0)}$ is finite-dimensional, the operator L on $\mathcal{H}^{(0)}$ is invertible [i.e., $L_1, L_2 \in \mathbb{C} \setminus \{0\}$]. Let $i \in \mathbb{N} \cup \{0\}$. If, for each $k = 0, \dots, i$, we have $\ker L_{11}^{(k)} = \{0\}$ ($L_{11}^{(k)}$ – restriction of $\Gamma_1^{(k)} L \Gamma_1^{(k)}$ to $\mathcal{E}^{(k)}$) [e.g., if $(L_1, L_2) \in (\mathbb{C}^+)^2$] then for all $k = 0, \dots, i$, the $L_*^{(k)} : \mathcal{U}^{(k)} \rightarrow \mathcal{U}^{(k)}$ is given by the previous representation. Furthermore, if $\mathcal{V}^{(k)} \cap \mathcal{J}^{(k)} = \{0\}$ for each $k = 0, \dots, i$, then the $Y_*^{(k)} : \mathcal{V}^{(k)} \rightarrow \mathcal{V}^{(k)}$ is given by the previous representation with L and $L_*^{(k)}$. Moreover, if

$$\mathcal{V}^{(k)} \cap \mathcal{E}^{(k)} = \{0\}, \mathcal{P}_j^{(k)} \cap \mathcal{U}^{(k)} = \{0\}, j = 1, 2, k = 0, \dots, i - 1$$

then for all $k = 0, \dots, i - 1$ we have

$$\dim \mathcal{U}^{(0)} = \dim \mathcal{U}^{(k)} = \dim \mathcal{V}^{(k)} = \dim \mathcal{U}^{(i)},$$

& $K^{(k)} : \mathcal{U}^{(k+1)} \rightarrow \mathcal{V}^{(k)}$, $Y_*^{(k)}$, $L_*^{(k+1)}$ are linked by the previous formulas.

Conclusion

- Conditions on materials (L_1, L_2) such that L on $\mathcal{H}^{(0)}$ and $L_{11}^{(k)}$ on $\mathcal{E}^{(k)}$ are invertible, i.e., $\ker L_{11}^{(k)} = \{0\}$ as $L_{11}^{(k)}$ is a Fredholm operator of index 0, is "easier" to check than $\mathcal{V}^{(k)} \cap \mathcal{J}^{(k)} = \{0\}$ or $\mathcal{P}_j^{(k)} \cap \mathcal{U}^{(k)} = \{0\}$, which is "highly dependent on the physical problem and structure of the composite" (e.g., Remark 47 in [2] & Sec. 19.4 in [3]).
 - If all conditions hold \Rightarrow Continued fraction expansion (!), infinite or terminating.
 - Conditions don't hold requires modification - future work.
- Continued fraction expansions - basis for bounding effective moduli of composites, see [3-6] (multiphase composites, see [4].)
- Future work - Extend all these results, i.e., field recursion method, to multiphase composites in abstract theory of composites.
- Open probs., e.g., [3, ch. 29]: "still an open (& interesting) question as to whether every matrix-valued analytic function $L_*(\lambda_1, \lambda_2, \dots, \lambda_n)$ satisfying the homogeneity, Herglotz, & normalization properties can be associated with a $(3, n)$ -subspace collection."

Auxiliary Slides

Quintessential example: conductivity equation

- For any $F \in \mathcal{H} = [L^2_{\text{per}}(\mathcal{D})]^d$,

$$\Gamma_0 F = \langle F \rangle, \quad \Gamma_1 F = [\nabla(\nabla \cdot \nabla)^{-1} \nabla \cdot] F, \quad \Gamma_2 F = [\nabla \times (\nabla \times \nabla \times)^{-1} \nabla \times] F,$$
$$\Lambda_j F = \chi_j F.$$

- Why? Helmholtz decomposition for periodic functions and

$$\nabla \cdot (\nabla \times) = 0$$

and

$$\begin{aligned} \nabla \times (\nabla \times \nabla \times)^{-1} \nabla \times &= \nabla \times (-\nabla \cdot \nabla)^{-1} \nabla \times \\ \nabla \times \nabla \times &= \nabla(\nabla \cdot) - \nabla \cdot \nabla \\ \nabla \times (\nabla) &= 0. \end{aligned}$$